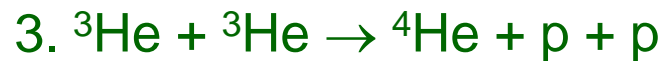
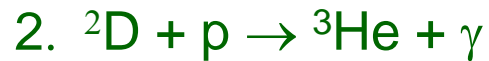
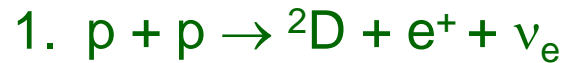


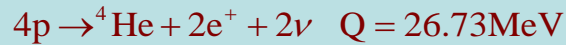
# Energy production in sun

- Sequence of nuclear reactions: proton-proton chain

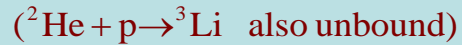


# Introduction

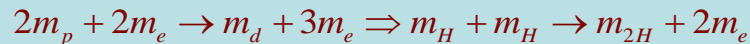
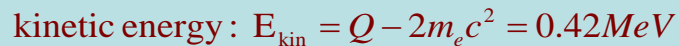
proton proton chain



possible first steps :



- Weak interaction process



kinetic energy is shared by positron and neutrino

- Theoretical cross section for the p + p reaction

$$\text{total Hamiltonian: } H = H_n + H_\beta$$

but  $H_\beta$  much smaller than  $H_n \Rightarrow$  perturbation theory

Fermis Golden Rule :

$$d\sigma = \frac{2\pi}{\hbar} \frac{\rho(E)}{v_i} |\langle f | H_\beta | i \rangle|^2$$

$\rho(E)$  statistical factor, density of final states

$v_i$  relative velocity in incident channel

$|\langle f | H_\beta | i \rangle|$  transition matrix element

$|i\rangle \equiv p + p$

$|f\rangle \equiv d + e^+ + \nu$

$$\rho(E) = \frac{dN}{dE} \text{ number of final states } dN \text{ in energy interval } dE \text{ (E and E + dE)}$$

$$\text{total phase space divided by unit volume} \quad dn = V \frac{4\pi p^2 dp}{h^3}$$

## Phase space

applying to pp reaction → states available for positrons and neutrinos

$$dN = dn_e \cdot dn_\nu = \left( V \frac{4\pi p_e^2 dp_e}{h^3} \right) \left( V \frac{4\pi p_\nu^2 dp_\nu}{h^3} \right)$$

assume : - zero rest mass of neutrino  $m_\nu = 0$

- zero recoil energy of deuteron  $E_{kin,d} = 0$

⇒ Energy is shared by positron and neutrino

$$E = E_e + E_\nu = E_e + cp_\nu$$

$$p_\nu = \frac{E}{c} - \frac{E_e}{c} \quad \text{for constant } p_e$$

$$\frac{dp_\nu}{dE} = \frac{1}{c} \quad dp_\nu = \frac{1}{c} dE \quad \frac{dp_\nu}{dE_e} = -\frac{1}{c}$$

$$\text{neutrino level density: } \frac{dn_\nu}{dE} = \left( V \frac{4\pi p_\nu^2 dp_\nu}{h^3} \right) = V \frac{4\pi}{h^3} p_\nu^2 \frac{dp_\nu}{dE} = V \frac{4\pi}{h^3} \frac{1}{c^2} (E - E_e)^2 \frac{1}{c}$$

$$dn_e = \left( V \frac{4\pi p_e^2 dp_e}{h^3} \right)$$

$$\rho(E) = \frac{dN}{dE} = dn_e \cdot \frac{dn_\nu}{dE} = \left( V \frac{4\pi p_e^2 dp_e}{h^3} \right) \cdot V \frac{4\pi}{h^3} \frac{1}{c^2} (E - E_e)^2 \frac{1}{c}$$

$$\rho(E) = V^2 \frac{16\pi^2}{h^6} \frac{1}{c^3} p_e^2 (E - E_e)^2 dp_e = \rho(E_e) dp_e$$

# Matrix element

Back to the cross section for the  $p + p$  reaction, cross section depends only on positron variables

Fermis Golden Rule :

$$d\sigma = \frac{2\pi}{\hbar} \frac{\rho(E_e)}{v_i} |\langle f | \mathbf{H}_\beta | i \rangle|^2 dp_e$$

$$\text{Matrix element : } H_{if} = \langle f | \mathbf{H}_\beta | i \rangle = \int_{Vol} \psi_f^* \mathbf{H}_\beta \psi_i d\tau = \int_{Vol} [\psi_d \psi_e \psi_\nu]^* \mathbf{H}_\beta \psi_i d\tau$$

Wave function of positron and neutrino : plane waves

$$\psi_e = \frac{1}{V^{1/2}} \exp(i\vec{k}_e \vec{r}) \quad \text{and} \quad \psi_\nu = \frac{1}{V^{1/2}} \exp(i\vec{k}_\nu \vec{r})$$

Leptons are normalized to the volume.

$$\text{For initial protons : } \int_{Vol} \psi_i^* \psi_i d\tau = 1$$

Deuteron wave function vanishes rapidly outside nuclear radius  $r_d = 1.7$  fm.

Mean neutrino energy from phase space :  $\bar{E}_\nu = 0.26$  MeV.

$$\Rightarrow k_\nu r_d = 2.2 \times 10^{-3}$$

$$\text{expansion of wave functions : } \psi_e = \frac{1}{V^{1/2}} [1 + i(\vec{k}_e \vec{r}) + \dots]$$

$$\psi_\nu = \frac{1}{V^{1/2}} [1 + i(\vec{k}_\nu \vec{r}) + \dots]$$

# Cross section

Matrix element after expansion :

$$H_{if} = \frac{1}{V} \int_{Vol} \Psi_d^* H_\beta \Psi_i d\tau \quad \text{strength of weak interaction } H_\beta \text{ is governed by coupling constant } g$$

(analog to electric charge as coupling constant  $\alpha$  in EM interaction)

$$H_{if} = \frac{1}{V} g \int_{Vol} \Psi_d^* \Psi_i d\tau = \frac{1}{V} g M_{space} M_{spin}$$

$$d\sigma = \frac{2\pi}{\hbar} \frac{\rho(E_e)}{v_i} |\langle f | H_\beta | i \rangle|^2 dp_e = \frac{2\pi}{\hbar} \frac{1}{v_i} V^2 \frac{16\pi^2}{h^6} \frac{1}{c^3} p_e^2 (E - E_e)^2 \left( \frac{1}{V} g M_{space} M_{spin} \right)^2 dp_e$$

$$\sigma = \frac{2\pi}{\hbar} \frac{1}{v_i} \frac{16\pi^2}{h^6} \frac{1}{c^3} g^2 M_{space}^2 M_{spin}^2 \int_0^E p_e^2 (E - E_e)^2 dp_e$$

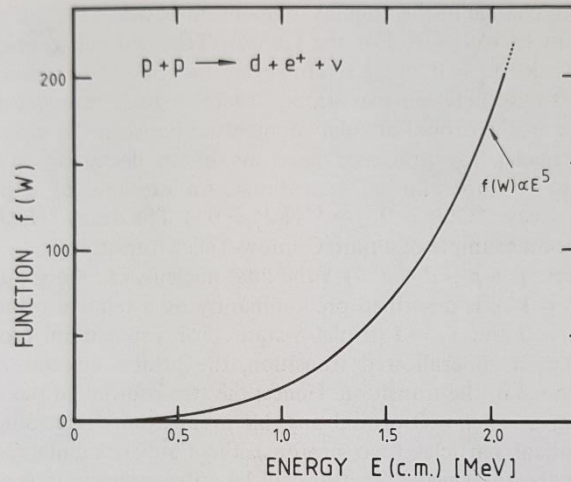
$$\text{new variable : } w = \frac{E + m_e c^2}{m_e c^2}$$

$$\text{substitution : } \int_0^E p_e^2 (E - E_e)^2 dp_e = \frac{(m_e c^2)^5}{c^3} \int_1^w (w_e^2 - 1)^{1/2} (w - w_e)^2 w_e dw_e = \frac{(m_e c^2)^5}{c^3} f(w)$$

integral is proportional to  $E^5$   $f(w) \propto w^5 \propto E^5$

$$\sigma = \frac{1}{v_i} \left[ \frac{2\pi}{\hbar} \frac{16\pi^2}{h^6} \frac{1}{c^3} \frac{(m_e c^2)^5}{c^3} \right] f(w) g^2 M_{space}^2 M_{spin}^2 \quad \text{constant values } \alpha = \left[ \frac{2\pi}{\hbar} \frac{16\pi^2}{h^6} \frac{1}{c^3} \frac{(m_e c^2)^5}{c^3} \right] = 1.45 \cdot 10^{70} \frac{1}{eV^2 \cdot s \cdot cm^6}$$

$$\sigma = \alpha \frac{1}{v_i} f(w) g^2 M_{space}^2 M_{spin}^2$$



The function  $f(W)$ , which is a measure of the total phase space available for the  $e^+ + \nu$  reaction, is shown as a function of energy  $E$ . As might be expected, the available phase space increases rapidly with  $E$ . For large values of energy, the phase space and hence the probability for the process increase as the fifth power of the energy.

An illustration of this function  $f(W)$  is presented in Figure 6.3. For values of  $W$ , i.e., large energies  $E$ , the function  $f(W)$  approaches the form  $f(W) \propto E^5$ .

The total cross section is given now by the equation

$$\frac{1}{v_i} \frac{2\pi}{\hbar} \frac{16\pi^2}{c^3 h^6} \frac{(m_e c^2)^5}{c^3} f(W) g^2 M_{\text{spin}}^2 M_{\text{space}}^2 \propto \frac{1}{v_i} f(W) g^2 M_{\text{spin}}^2 M_{\text{space}}^2, \quad (6)$$

where  $g$  represents the constants, which, when replaced by their numerical values, result in

$$\frac{m^5 c^4}{2\pi^3 \hbar^7} = 1.45 \times 10^{70} \text{ eV}^{-2} \text{ s}^{-1} \text{ cm}^{-6}.$$

Observations of allowed weak-interaction phenomena indicate (Bla6) that among all possible types of couplings (vector, scalar, tensor, axial vector, and axial scalar) the vector and axial vector couplings are the most common.

# Weak interaction couplings

(i) Spins of leptons in final state are  $\uparrow\downarrow$  antiparallel

$$s_e + s_\nu = s_{\text{tot}} = 0$$

Fermi or vector coupling  $g=C_V$

change in angular momentum:  $\Delta J=0$

e.g. Fermi transition:  $^{14}\text{O}(J^\pi_i=0^+) \rightarrow ^{14}\text{N}(J^\pi_f=0^+)$

(ii) Spins of leptons in final state are  $\uparrow\uparrow$  parallel

$$s_e + s_\nu = s_{\text{tot}} = 1$$

Gamow-Teller or axial-vector coupling  $g=C_A$

change in angular momentum:  $\Delta J=0, \pm 1$

*However,  $\Delta J=0$  not possible between two states of zero angular momentum.*

e.g. GT transition:  $^6\text{He}(J^\pi_i=0^+) \rightarrow ^6\text{Li}(J^\pi_f=1^+)$

# Implications for pp reaction



- (i) Spins of ground state of deuteron, final state  ${}^2\text{H}$  ( $J^\pi_f=1^+$ ) described by  $l_f=0$  and  $s_f=1$ , triplet s-state, only small d-state admixture.
- (ii) Max. probability to react is related to process without change in angular momentum. Two protons are in  $l_i=0$ , implies a symmetric wave function for  $\Psi_{\text{space}}$ .
- (iii) Two identical Fermions (2 protons)  $\rightarrow$  Pauli principle requires singlet configuration  $s_i = 0$  ( $\uparrow\downarrow$ ) to get an antisymmetric wave function. ( $\Psi_{\text{tot}} = \Psi_{\text{space}} \Psi_{\text{spin}}$ ).
- (iv)  $s_i=0, l_i=0 \rightarrow s_f=1, l_f=0$   
Gamow-Teller-Transition with coupling constant  $g=C_A$
- (v)  $C_A$  is determined from lifetimes of GT transitions in neutron, triton,  ${}^6\text{He}$  decay.



# Final results for pp reaction



(vi)  $M_{\text{spin}}$

summing over final states,  
averaging over initial states,  
division by 2 due to the  
fact of identical particles

$$M_{\text{spin}} = 3/2$$

(vii)  $M_{\text{space}}$

numerical integration  
of product of radial wave  
functions of protons and  
deuteron.

$$M_{\text{space}} = \int_0^{\infty} \chi_f(r) \chi_i(r) r^2 dr$$

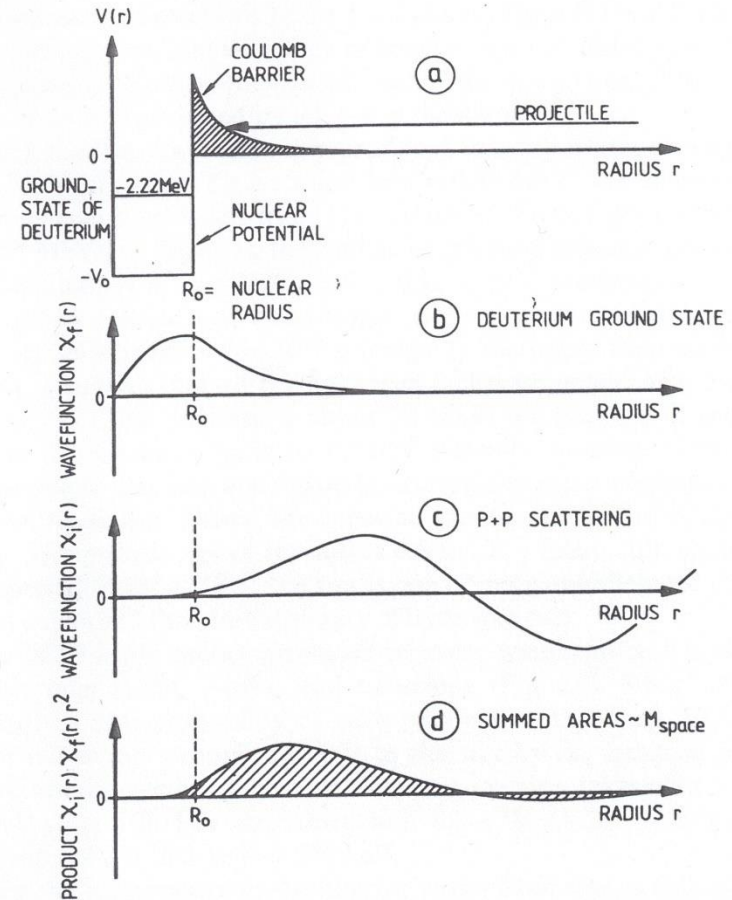


FIGURE 6.4. Shown schematically are a few ingredients used in the numerical evaluation of the space matrix element  $M_{\text{space}}$  for the  $p + p \rightarrow d + e^+ + \nu$  reaction. The potential is shown in (a), where, for a given nuclear radius  $R_0$ , the observed binding energy of the deuteron determines the potential depth  $V_0$ . The deuteron radial wave function  $\chi_f(r)$  is determined by the potential  $V(r)$ . Because of the loosely bound ground state,  $\chi_f(r)$  extends far outside  $R_0$  with appreciable amplitudes (b). The initial wave function  $\chi_i(r)$  is obtained from  $p + p$  elastic scattering data, which gives (c) a small amplitude for  $r \leq R_0$  and has the usual oscillating pattern of a plane wave for  $r \gg R_0$ . The radial integrand in  $M_{\text{space}}$  (d) then has its major contributions in regions far outside  $R_0$  (hatched areas).

# Final results for pp reaction



(vi)  $M_{\text{spin}}$  summing over the final states, averaging over the initial states, dividing by 2 due to the fact of identical particles  $M_{\text{spin}} = 3/2$ .

(vii)  $M_{\text{space}}$  numerical integration of product of radial wave functions of protons and deuteron.

$$M_{\text{space}} = \int_0^{\infty} \chi_f(r) \chi_i(r) r^2 dr$$

Final result: Total cross section at  $E_p=1\text{MeV}$

$$\sigma(E) \sim 10^{-47} \text{ cm}^2$$

non resonant reaction, smooth variation with energy at  $T_6=15$

$$\langle \sigma v \rangle_{pp} = 1.19 \times 10^{-43} \left[ \text{cm}^3 \frac{1}{\text{s}} \right]$$

# What is happening inside the sun?



Sun interior is a mixture of H and He

Assumption:  $X_H \sim X_{He} \sim 0.5$       density  $\rho \sim 100 \text{ g/cm}^3$

Mean lifetime of proton against combustion

$$\tau_H = \frac{1}{N_H \langle \sigma v \rangle_{pp}} = 0.9 \cdot 10^{10} \text{ year}$$

Overall rate of the conversion of four protons into  $^4\text{He}$  is governed by the pp reaction!

$\tau \sim 10^{10}$  years  $\sim$  age of oldest known stars!

# What is happening inside the sun next?



Burning of deuteron:



$$\frac{dd}{dt} = r_{pp} - r_{pd} \quad r_{12} = \frac{N_1 N_2}{1 + \delta_{12}} \langle \sigma v \rangle_{12}$$

$$\text{quasi equilibrium: } \frac{dd}{dt} = 0$$

$$\left( \frac{d}{p} \right) = \frac{\langle \sigma v \rangle_{pp}}{2 \langle \sigma v \rangle_{pd}} \quad \left( \frac{d}{p} \right) = 5.6 \times 10^{-18} \quad \text{at } T_6 = 5$$

Mean lifetime of deuteron against combustion  $\tau(\text{D}) = 1.6 \text{ s}$  !

# Reaction rates and network

$$\left(\frac{dN_1}{dt}\right)_2 = -(1 + \delta_{12})r = -N_1N_2\langle\sigma v\rangle$$

$$= -N_1\rho N_A \frac{X_2}{A_2}\langle\sigma v\rangle = -N_1\rho N_A Y_2\langle\sigma v\rangle$$

$$r = \frac{1}{1 + \delta_{12}} N_1N_2\langle\sigma v\rangle$$

$$N_i = \rho N_A \frac{X_i}{A_i} = \rho N_A Y_i$$

Number density  $N$

Matter density  $\rho$

Avogadro's constant  $N_A$

Mass fraction  $X$

Atomic mass  $A$

Mol fraction  $Y$

The solar ppl chain consists of the following reactions:  
 $p(p, e^+ \nu) d$  (11),  $d(p, \gamma) {}^3\text{He}$  (12),  ${}^3\text{He}({}^3\text{He}, 2p) {}^4\text{He}$  (33)  
with  $Y_1 = Y_p$ ,  $Y_2 = Y_d$ ,  $Y_3 = Y_{{}^3\text{He}}$ ,  $Y_4 = Y_{{}^4\text{He}}$

$$\dot{Y}_1 = \rho N_A (-\langle\sigma v\rangle_{11} Y_1^2 - \langle\sigma v\rangle_{12} Y_1 Y_2 + \langle\sigma v\rangle_{33} Y_3^2)$$

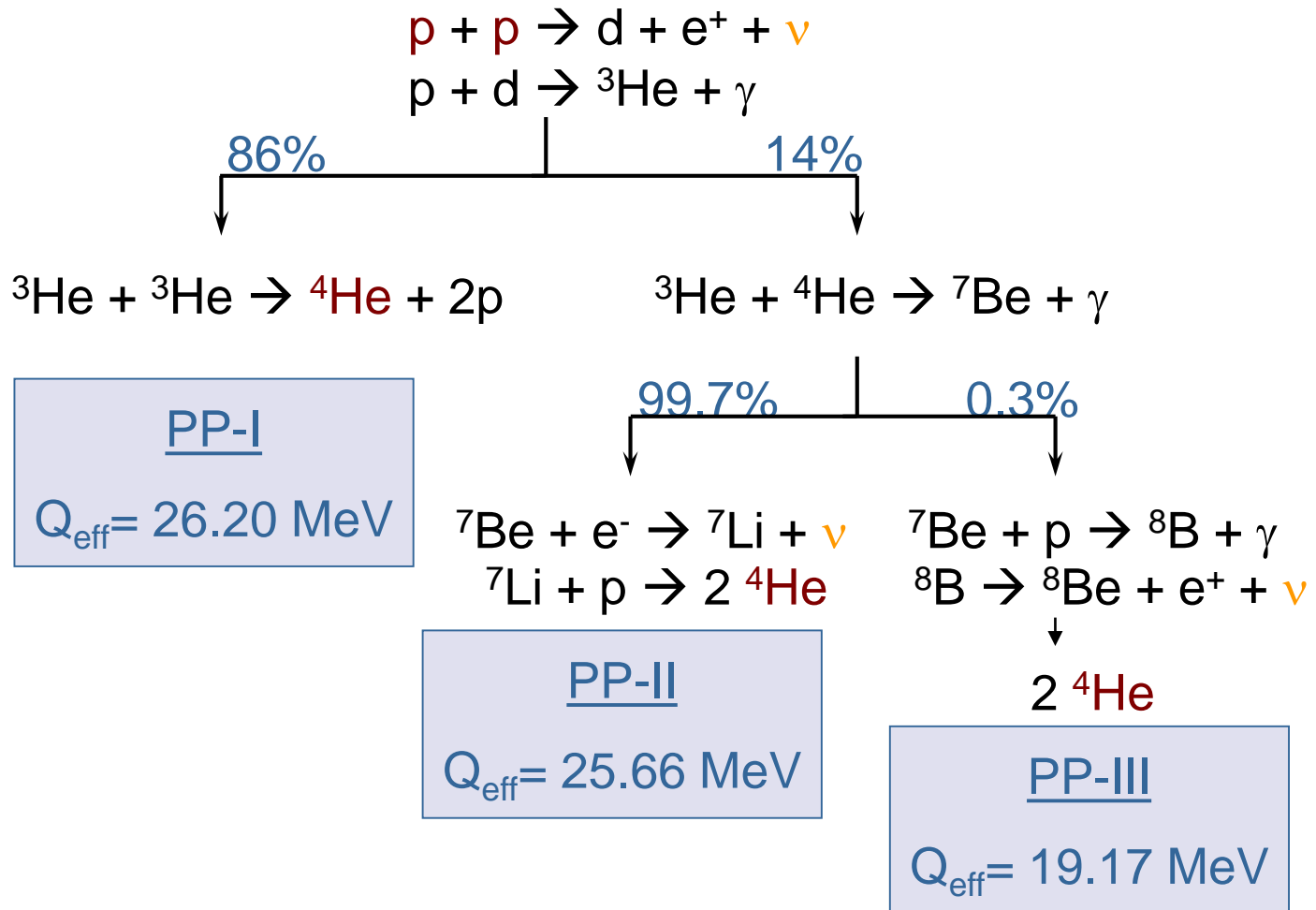
$$\dot{Y}_2 = \rho N_A (+\frac{1}{2}\langle\sigma v\rangle_{11} Y_1^2 - \langle\sigma v\rangle_{12} Y_1 Y_2)$$

$$\dot{Y}_3 = \rho N_A (+\langle\sigma v\rangle_{12} Y_1 Y_2 - \langle\sigma v\rangle_{33} Y_3^2)$$

$$\dot{Y}_4 = \rho N_A (+\frac{1}{2}\langle\sigma v\rangle_{33} Y_3^2)$$

The factors 1/2 avoid double-counting of identical nuclei.

# proton-proton chain



net result:  $4p \rightarrow {}^4\text{He} + 2e^+ + 2\nu + Q_{\text{eff}}$